

## APPLICATION OF VARIATION METHODS FOR CALCULATION ELASTOMERIC ELEMENTS OF A DIFFICULT CONFIGURATION

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**Abstract:** *For an efficient utilization of elastomeric elements in modern mechanical engineering, including vibration insulation, it is necessary to be able to calculate the characteristics of compressive stiffness of elastomeric elements. Many elastomeric elements have the geometrical form of a complex configuration. In this work the design procedure the compressive stiffness characteristics of type "force – settlement", for elastomeric elements of any geometrical form is offered at static loading. . In the given work for reception of the approached decisions it is offered to use direct methods using functional considered by V.Prager [1] and converted to weakly compressible materials. Methodology use is shown on an example of calculation of the concrete rubber absorber.*

*Key words: compressive stiffness, rubber, shock-absorber, subareas.*

### 1. PROBLEM DESCRIPTION

For an efficient utilization of elastomeric elements in modern mechanical engineering, including vibration insulation, it is necessary to be able to calculate the characteristics of compressive stiffness of elastomeric elements. Many elastomeric elements have a complex geometric configuration, which does not permit the application of variational methods of calculation [2, 3], effectively used in the calculation of elastomeric elements simple geometry. But for the elastomeric elements of complex configuration or consisting of well contacting parts made of various

materials, using functionals [2, 3] virtually impossible to choose a continuous function of displacement required to satisfy the geometric conditions and the conditions of continuity throughout the volume of complex elastomeric element. In this paper considered the method of calculating the rigidity characteristics under static load, type of force - settlement, for the elastomeric elements of any geometric shapes. Deformations are assumed to be small. It is proposed to use the variational principle proposed V.Prager [1] using discontinuous displacement function and modified for weakly compressible and incompressible materials.

### 2. DECOMPOSITION OF THE VOLUME OF THE PRODUCTS ON A SUBAREAS

Let's consider the elastomeric element with the complex geometrical form with volume  $V$  and the area of a surface  $F$ :

$$V = \sum_{n=1}^N V_n \quad F = \sum_{n=1}^N F_n \quad (1)$$

where:

$V_n$  - one-coherent regular subareas;  
 $N$  - number of subareas, received as a result of crushing elastomeric element;  
 $F$  - total surface area of all subdomains elastomeric element.

The surface limiting  $n$ -th subarea, looks like

$$F_n = F_\sigma^n + F_u^n + \Gamma_n \quad (2)$$

where:

$F_\sigma^n$  – loading surface;

$F_u^n$  – surface of the fixing;  
 $\Gamma_n$  – surface area of contact subarea  
partitioning elastomeric element.

On the crushing surface  $\Gamma_n$  must be  
satisfied the conditions of continuity:  
- of displacements:

$$u_i^n = u_i^{n+1} \quad (3)$$

- and of stresses:

$$\sigma_{ij}^n m_j^n = -\sigma_{ij}^{n+1} m_j^{n+1} \quad (4)$$

Where the index „n” specifies a current  
issue of a subarea of crushing and  $m_j^n$  and  
 $m_j^{n+1}$  - directing cosines normals external,  
accordingly to  $V_n$  and  $V_{n+1}$  on  $\Gamma_n$  and  $\Gamma_{n+1}$ .  
If you use only the external normal to  $V_n$ ,  
which at  $\Gamma_{n+1}$  for  $V_{n+1}$  is internal, then  
 $m_j^{n+1} = -m_j^n$  and the condition (4) reads:

$$\sigma_{ij}^n m_j^n = \sigma_{ij}^{n+1} m_j^n \quad (5)$$

Hereinafter, for brevity, over repeated  
lower indices are summed, and the comma  
denotes a partial derivative.

Contact surface  $\Gamma_n$  may be artificial in the  
geometric decomposition of the  
elastomeric element, or natural, if the  
physical and mechanical characteristics of  
the material ( $G, \mu$ ) at the contact surface  
change abruptly, that is, a volume  $V$   
composed of different materials.

Partition of the elastomeric element in the  
subareas extends the permissible class of  
unknown functions in the piecewise  
smooth and piecewise continuous functions  
with piecewise smooth and piecewise  
continuous derivatives [1, 4]. Discontinuity  
of displacements and efforts on the surface  
 $\Gamma_n$  denote:

$$\begin{aligned} u_i^n - u_i^{n+1} &= \{u_i^n\}, \\ \sigma_{ij}^n m_j^n - \sigma_{ij}^{n+1} m_j^n &= \{\sigma_{ij}^n m_j^n\} \end{aligned} \quad (6)$$

If the surface areas  $\Gamma_n$  of subareas  $V_n$   
displacement components and efforts (all

or only some) does not satisfy the  
conditions of continuity (3) and (4), then,  
using the designation (6), we can write:

$$\{\sigma_{ij}^n m_j^n u_i^n\} = \{\sigma_{ij}^n m_j^n\}^I u_i^n + \sigma_{ij}^n m_j^n \{u_i^n\}^{II} \quad (7)$$

where the indices I and II, respectively,  
indicate that the summation extends only to  
the displacement components and efforts  
that do not satisfy the conditions on the  
surface  $\Gamma_n$ .

Suppose that in each subareas partitioning  
the have the required properties of  
continuity and differentiability.

Variational principle for discontinuous  
functions during the fragmentation of the  
field on a subareas is given in [1]. Using  
the method of Lagrange multipliers, can be  
generalized to the case of weakly  
compressible and incompressible materials,  
which include majority of elastomeric  
materials.

### 3. MATHEMATICAL MODEL

When solving boundary value problems of  
static elasticity theory for incompressible  
and weakly compressible materials easier  
as the unknown functions to choose  
displacement  $u_i$  and the function of  
hydrostatic pressure  $s$ , which, for small  
strains, leads to a mathematical model [2,3]:  
Equation of equilibrium:

$$G \left[ \nabla^2 u_i + \frac{3}{2(1+\mu)} s_i \right] + f_i = 0 \quad (8)$$

Volumetric deformation:

$$u_{j,j} = \frac{3(1-2\mu)}{2(1+\mu)} s \quad (9)$$

Deformations:

$$\varepsilon_{ij} = 0,5 (u_{i,j} + u_{j,i}) \quad (10)$$

Stress

$$\sigma_{ij} = G [u_{i,j} + u_{j,i} + 3\mu / (1+\mu) s \delta_{ij}] \quad (11)$$

Forces boundary conditions:

$$G[u_{i,j} + u_{j,i} + 3\mu / (1+\mu) s \delta_{ij}] n_j = P_i \quad (12)$$

Displacements boundary conditions:

$$u_i = u_{0i} \quad (13)$$

When partition the element into a subarea to a mathematical model (8) - (13) to add conditions for docking (3) and (4).

In determining the integral characteristics, of type “force – settlement”, of the elastomeric element boundary problem (8) - (13) (without crushing the elastomeric element in the subareas) easier to solve a variational method using the Ritz procedure for the functional [2, 3]:

$$J(u_i^n, s^n) = G \int_V \left[ \frac{1}{2} (u_{i,j}^n u_{j,i}^n + u_{i,j}^n u_{j,i}^n) + \frac{3\mu}{1+\mu} s^n u_{i,i}^n - \frac{9(1-2\mu_n)}{4(1+\mu_n)^2} s^{n2} \right] dV \quad (14)$$

Using the variational principle of V.Prager [1] and applying the method of undetermined Lagrange multipliers with a functional (14) be the boundary value problem (8) - (13) with the conditions of the joining (3) and (4) replaced by the variational problem with discontinuous function of demand on the surfaces of the partition  $\Gamma_n$  for the functional:

$$J^*(u_i, s) = \sum_{n=1}^N J(u_i^n, s^n) - \sum_{n=1}^{N-1} G_n \int_{\Gamma_n} [(u_{i,j}^n + u_{j+i}^n) + \frac{3\mu_n}{1+\mu_n} s^n \delta_{i,j} m_j \{u_i^n\}'' ] d\Gamma_n \quad (15)$$

where in each subare:

$u_i^n$  – displacement;

$s^n$  - function of hydrostatic pressure;

$G_n$  - modulus of elasticity in shear;

$\mu_n$  - Poisson's ratio.

Using the functional (15), the choice of displacement functions  $u_i^n$  only need to follow geometric boundary conditions (13), as a function of hydrostatic pressure  $s_n$  in each subarea can be selected

independently, not caring about its continuity at the boundary of the partition  $\Gamma_n$ .

As an example, consider obtaining dependence “force – settlement” to the hollow cylindrical absorber with core under axial compression (Fig. 1).

Elastomeric layer is still sealed with absolutely rigid top and bottom plates and a non-deformable inner core. For this shock absorber is difficult to write expressions for the displacement common to all elastomeric layer.

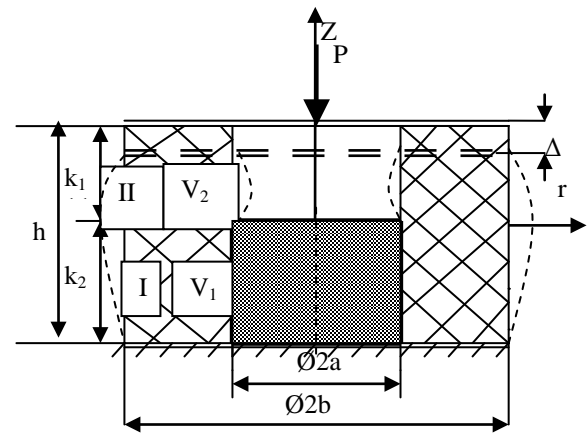


Fig. 1. The hollow cylindrical absorber with the arrester

The considered absorber is broken on two parts on border of a thrust block in parallel a shaft or. At use functional  $J^*(u_i, s)$  (15), choosing  $u_i^n$  and  $s_n$ , it is enough to satisfy to geometrical boundary conditions on external surface  $F_u$ . Shock absorber is divided into two subareas (see Figure 1). All functions with an index „1” it is carried to a subarea I, and with an index „2” to a subarea II. We believe that the geometry of the elastomeric layer can not take into account the compressibility of the elastomer, that is, believe that the Poisson coefficient  $\mu = 0,5$ .

The main boundary conditions will be:

$$\begin{aligned} u_1(r, -k_2) = u_2(r, k_1) &= 0 \\ u_1(a, z)_{|0 \leq z \leq -k_2} = w_1(a, z)_{|0 \leq z \leq -k_2} &= 0 \\ w_1(r, -k_2) &= 0 \\ w_2(r, k_1) &= -\Delta \end{aligned} \quad (16)$$

where: the functions  $u_i$  and  $w_i$  - displacement, respectively, on the axis of  $r$  and  $z$ .

Dependence “force – settlement” it is defined from the equation of balance of the top base of the absorber:

$$2\pi \int_a^b \sigma_{2zz}|_{z=k_1} r dr = -P \quad (17)$$

On a surface of splitting of a condition of ideal contact piece look like:

$$\begin{aligned} u_1(r,0) &= u_2(r,0) \\ w_1(r,0) &= w_2(r,0) \\ \sigma_{1rz}(r,0) &= \sigma_{2rz}(r,0) \\ \sigma_{1zz}(r,0) &= \sigma_{2zz}(r,0) \end{aligned} \quad (18)$$

We choose conveyances  $u_n$ ,  $w_n$  and function  $s_n$  whenever possible in the most simple kind with the account only conditions (11) and prospective character of deformation.

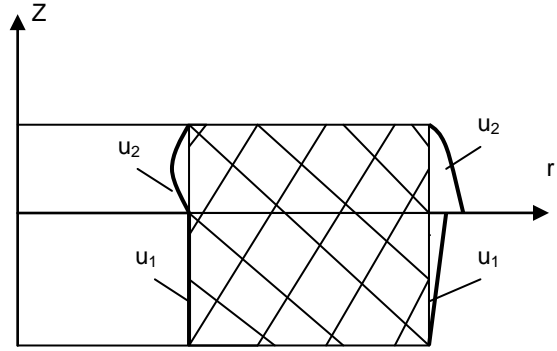
$$\begin{aligned} u_2 &= B_1 r z (z - k_1) + B_2 (r - a)(z - k_1) \\ w_2 &= -\frac{\Delta z}{k_1} + A_1 (r - a)(z - k_1) \\ s_2 &= C_2 \\ u_1 &= B_3 (r - a)(z + k_2) \\ w_1 &= A_2 (r - a)(z + k_2) \\ s_1 &= C_1 \end{aligned} \quad (19)$$

where

$A_1, A_2, B_1, B_2, C_1, C_2, \Delta$  – unknown constants.

Functions (19) have on section height approximately following appearance for conveyances at  $r = a$ ;  $b$ :

for  $u$



for  $w$

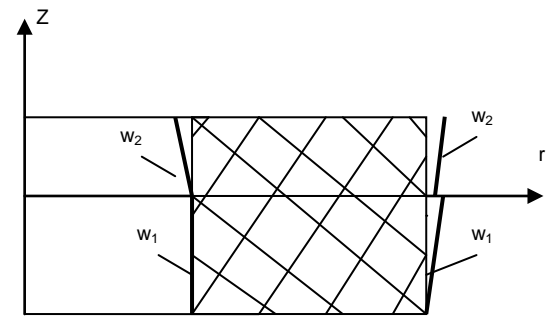


Fig. 2. Expected character of deformation

After integration it is received that functional  $J^*$  depends only on unknown constants:

$$J^* = J(A_1, A_2, B_1, B_2, B_3, C_1, C_2, \Delta) \quad (20)$$

From a condition stationarity

$$\frac{\partial J^*}{\partial(A_1, A_2, B_1, B_2, B_3, C_1, C_2, \Delta)} = 0 \quad (21)$$

We receive system of the algebraic equations. For dependence “force – settlement”:

$$\Delta = \frac{P k_1}{2\pi G a^2 (1 - \alpha^2)} \frac{D_1}{D}, \quad \alpha = \frac{a}{b} \quad (22)$$

where:

$D, D_1$  - determinants of algebraic equations (21), an expression which, due to the complexity of writing, are not given.

For absorber:  $h = 3$  cm,  $a = 2$  cm,  $b = 4,1$  cm,  $k_2 = 1,4$  cm,  $G = 7$  kg/cm<sup>2</sup>, in Figure 3 shows the results: “1” line - the formula (22) (using MatCad); “2” line - experiment.

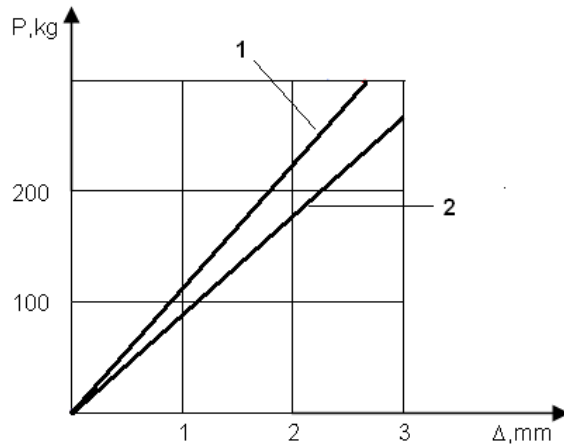


Fig. 3. Dependence force - settlement for the hollow cylindrical absorber with the arrester

If to accept that  $k_2 = 0$  then from expressions (21) and (22) we will receive the absorber without an internal thrust block.

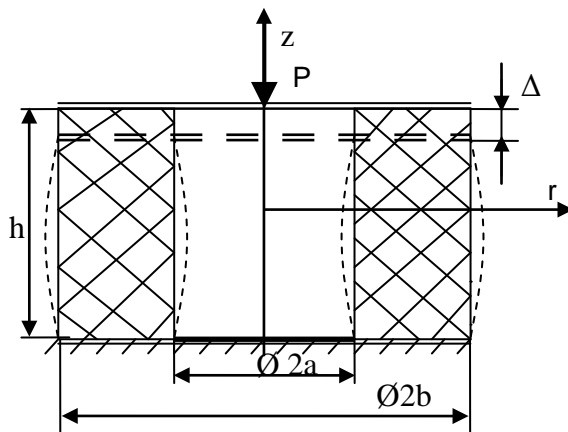


Fig. 4. The hollow cylindrical absorber without the arrester

In this case the dependence of the force - settlement will have the form:

$$\Delta = \frac{Ph}{2\pi Gb^2(1-\alpha^2)} D_3^{-1} \quad (23)$$

where

$$D_3 = 1,5 + 0,62\rho^2(1+\alpha^2) \left[ 1 + \frac{\rho^2(1-\alpha^2)\alpha^2}{0,26(1-\alpha^2) - \rho^2\alpha^2 \ln\alpha} \right]$$

$$\alpha = \frac{a}{b}; \quad \rho = \frac{b}{h}$$

Figure 5 The predicted results are: "2" line - according to the formula (23); "1" line - solution obtained in work [2].

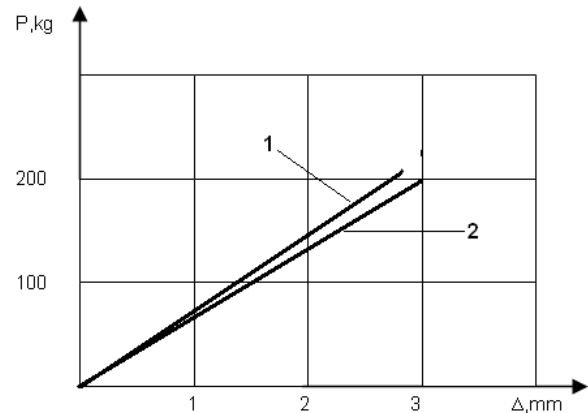


Fig. 5. Dependence force - a deposit for the hollow cylindrical absorber without the arrester

#### 4. CONCLUSION

The proposed method allows in case of partition of the investigated area  $V$  on subarea  $V_n$ , using functional (15) and discontinuous functions sought to obtain integrated dependences of type "force - settlement".

#### 5. ACKNOWLEDGEMENT

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#### 6. REFERENCES

1. Pragyer V. Variatsionniye printsipi linyeynoy staticheskoj tyeorii ooproogosti pri razrivnih smyeshshyenyah, dyeformatsiyah i napryazhyenyah. Myehanika. Sb.pyeryevodov.M., № 5, 1970.
2. Lavyendyel E.E. Raschyeti ryezino-tyehnichyeskikh izdyeliy. M., 1976. 230 s
3. Gontsa V.F. Vliyaniye slaboy szhimayemosti na ryeshyeniye zadach tyeorii ooproogosti dlya nyeszhimayemogo matyeriala. «Voprosi dinamiki i

prochnosti». Riga, 1970, vip. 20, s. 185-189.

4. Koorant R., Gil'byert D. Myetodi matyematicheskoj fiziki. T.I, II, M.-L.,1951.

5. Euler, M., Beigholdt, H.-A. Ermittlung von Kriechfunktionen für das viskoelastische Materialverhalten von Holf im Zugversuch / LACER No.4, Universitat Leipzig, 1999 -p.319-334.

6. Lyeykand, N.A., Lavyendyel, E.E., Goryelik, B.M. i dr. Ekspyrimyental'noye isslyedovaniye konstant ooproogosti ryezina 2959 i 51-1673. - V kn.: Vsyenaooch.-tyehn. konfyeryentsiya po myetodikye raschyeta izdyeliy iz visokoelastichnih matyerialov. Riga, 1980, s. 142 - 143.

7. Wood L.A. Values of physical constants of different rubbers.// In: Smith physical tabeles. 9th rev. Ed. Wash., 1956, p. 234-235.

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